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T. Fehlinger, W. Freeden, S. Gramsch,
C. Mayer, D. Michel, M. Schreiner

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from (Geostrophic) Ocean Flow**

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Local Modelling of Sea Surface Topography from (Geostrophic) Ocean Flow

T. Fehlinger¹, W. Freeden^{1*}, S. Gramsch¹, C. Mayer¹, D. Michel², and M. Schreiner³

¹ University of Kaiserslautern
Geomathematics Group
67663 Kaiserslautern
P.O. Box 3049
Germany

² Proudman Oceanographic Laboratory
6 Brownlow Street
Liverpool L3 5DA
UK

³ University of Buchs NTB
Laboratory for Industrial Mathematics
Werdenbergstrasse 4
CH-9471 Buchs
Switzerland

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Abstract: This paper deals with the problem of determining the sea surface topography from geostrophic flow of ocean currents on local domains of the spherical Earth. In mathematical context the problem amounts to the solution of a spherical differential equation relating the surface curl gradient of a scalar field (sea surface topography) to a surface divergence-free vector field (geostrophic ocean flow).

At first, a continuous solution theory is presented in the framework of an integral formula involving Green's function of the spherical Beltrami operator. Different criteria derived from spherical vector analysis are given to investigate uniqueness. Second, for practical applications Green's function is replaced by a regularized counterpart. The solution is obtained by a convolution of the flow field with a scaled version of the regularized Green function.

Calculating locally without boundary correction would lead to errors near the boundary. To avoid these Gibbs phenomena we additionally consider the boundary integral of the corresponding region on the sphere which occurs in the integral formula of the solution. For reasons of simplicity we discuss a spherical cap first, that means we consider a continuously differentiable (regular) boundary curve. In a second step we concentrate on a more complicated domain with a non continuously differentiable boundary curve, namely a rectangular region. It will turn out that the boundary integral provides a major part for stabilizing and reconstructing the approximation of the solution in our multiscale procedure.

1 Introduction

As a crucial factor for weather and climate, the oceans are of immediate importance. By modelling oceanic currents on the sphere (and in a first approximation the regions under consideration may be assumed to be part of a spherically shaped Earth), we gain a better understanding of meteorological processes. The point of departure for our intention to determine the sea surface topography is the geostrophic flow equation derived in conventional form from the basic hydrodynamic equation (see, e.g., [1, 5, 11, 12]). As a scalar field on the spherical Earth, the sea surface topography H consists of two ingredients. On the one hand, on an Earth at rest, the water masses would align along the geoid (given as the deviation from

* Corresponding author, e-mail: freeden@mathematik.uni-kl.de, Phone: +49 631 205 3867, Fax: +49 631 205 4736

the standard reference ellipsoid used in physical geodesy (see [7, 8])). On the other hand, satellite measurements provide altimetric data of the actual sea surface height which in marine geodesy are also used (in relation to the standard reference ellipsoid). The difference between these quantities can, of course, be considered as the actual sea surface topography. In other words, the sea surface topography H is defined as the deviation of the ocean surface from the geoidal surface which is here assumed to be due to the geostrophic component of the ocean currents. The surface geostrophic velocity of the ocean currents then reads as follows (see [1, 5, 11, 12])

$$L_\xi^* H(\xi) = h(\xi), \quad \xi \in \Omega^{oc}, \quad (1.1)$$

where Ω^{oc} is the subset of the Earth's sphere $\Omega \subset \mathbb{R}^3$ being covered by oceans (for simplicity, we associate Ω with the unit sphere in \mathbb{R}^3), L_ξ^* is the surface curl gradient acting on $\xi \in \Omega^{oc}$, defined by $L_\xi^* = \xi \wedge \nabla_\xi^*$ with ∇_ξ^* the surface gradient. Further, $h(\xi)$ is the product of three factors, namely (i) the quotient between the mean equatorial Earth's radius R and the acceleration of gravity $g(\xi)$ at ξ , (ii) the Coriolis force coefficient $c(\xi) = 2|\omega|(\varepsilon^3 \cdot \xi)$ (with $|\omega|$ the angular velocity of the Earth), and (iii) the velocity field $v(\xi)$ at ξ . In detail,

$$h(\xi) = \frac{2R}{g(\xi)} |\omega| (\varepsilon^3 \cdot \xi) v(\xi), \quad \xi \in \Omega^{oc}. \quad (1.2)$$

Clearly, for all $\xi \in \Omega^{oc}$, the geostrophic flow h given by (1.1) is perpendicular to the tangential surface gradient $\nabla^* H$ of the sea surface topography on Ω^{oc} . This is a remarkable feature of the geostrophic velocity field. The currents flow along and not across the lines of constant sea surface topography.

In this paper we are concerned with the so called *Topography Problem* of geostrophic ocean circulation in a normal region $\Gamma \subset \Omega^{oc}$: Let the vectorial velocity field h be known for a finite subset of points $\{\eta_1, \dots, \eta_N\}$ on Γ , and the scalar sea surface topography H be known for a set $\{\tilde{\eta}_1, \dots, \tilde{\eta}_{\tilde{N}}\}$ on the boundary $\partial\Gamma$ of Γ . Find an approximation of H from the discrete data $\{(\eta_i, h(\eta_i))\}_{i=1, \dots, N}$ and $\{(\tilde{\eta}_i, H(\tilde{\eta}_i))\}_{i=1, \dots, \tilde{N}}$ on the whole domain $\bar{\Gamma} = \Gamma \cup \partial\Gamma$.

In oceanography the sea surface topography is conventionally represented (on the whole Earth's surface) on the sphere Ω by a conventional Fourier (orthogonal) expansion in terms of spherical harmonics $Y_{n,j}$ (see, e.g., [10]) thereby assuming $\int_\Omega H(\xi) Y_{0,1}(\xi) d\omega(\xi) = \int_{\Omega^{oc}} H(\xi) Y_{0,1}(\xi) d\omega(\xi) = 0$ ($d\omega$ denotes the surface element). This approach ignores the local formulation of the problem (1.1) for the oceanic part $\bar{\Omega}^{oc}$ of the (spherical) Earth's surface. Instead it leads to an orthogonal series expansion of the velocity field in terms of surface divergence free vector spherical harmonics $L^* Y_{n,j}$ on the whole (spherical) Earth. But obviously, these vector types of polynomial functions are far from being suitable for purposes of approximation on the oceanic area Ω^{oc} . First, the constituting ingredients of spherical harmonics show certain phenomena of (global) periodicity at least when the classical basis system involving associated Legendre functions is used. In consequence, the assumptions on continents influence the approximation on oceans. Second, boundary effects along the coast lines like the Gibbs phenomenon are not avoidable by use of spherical harmonics, i.e., by use of non-space-localizing polynomials. So it is really necessary for geoscientific practice to develop on a geostrophic theory of ocean circulation in restriction to the oceanic parts Ω^{oc} on the Earth's surface thereby using specific space-localizing (trial) kernel vector functions for purposes of approximation on Ω^{oc} . These requirements will be investigated in the approach presented below.

2 Preliminaries

Throughout this paper we need a number of differential operators on the unit sphere $\Omega \subset \mathbb{R}^3$ which are listed in Table 1 (see, e.g., [4]).

Table 1 Differential operators

Symbol	Differential Operator
∇_x	Gradient at x
$\Delta_x = \nabla_x \cdot \nabla_x$	Laplace operator at x
∇_ξ^*	Surface gradient on the unit sphere Ω at ξ
$L_\xi^* = \xi \wedge \nabla_\xi^*$	Surface curl gradient on the unit sphere Ω at ξ
$\Delta_\xi^* = \nabla_\xi^* \cdot \nabla_\xi^* = L_\xi^* \cdot L_\xi^*$	Beltrami operator on the unit sphere Ω at ξ
$\nabla_\xi^* \cdot$	Surface divergence on the unit sphere Ω at ξ
$L_\xi^* \cdot$	Surface curl on the unit sphere Ω at ξ

It should be noted that the operators ∇^* , L^* , and Δ^* will be always used in a coordinate-free representation, thereby avoiding any kind of singularities at the poles. Moreover, following the nomenclature of [4] we denote by ∇^* the surface divergence on Ω and by L^* the surface curl on Ω . Clearly, $\Delta^* = L^* \cdot L^* = \nabla^* \cdot \nabla^*$. Note that the operators ∇^* , L^* , Δ^* show special features in certain situations (for more details the reader is referred to [4]). For example, let $\eta \in \Omega$ be fixed. If G is of class $C^{(1)}[-1, 1]$ and $G' \in C^{(0)}[-1, 1]$ denotes its derivative, then we find

$$\begin{aligned} \nabla_\xi^* G(\xi \cdot \eta) &= G'(\xi \cdot \eta)(\eta - (\eta \cdot \xi)\xi), \quad \xi \in \Omega, \\ L_\xi^* G(\xi \cdot \eta) &= G'(\xi \cdot \eta)(\xi \wedge \eta), \quad \xi \in \Omega, \end{aligned} \quad (2.1)$$

whereas for $G \in C^{(2)}[-1, 1]$

$$\Delta_\xi^* G(\xi \cdot \eta) = (L^* \cdot L^*)G(\xi \cdot \eta) = -2(\xi \cdot \eta)G'(\xi \cdot \eta) + (1 - (\xi \cdot \eta)^2)G''(\xi \cdot \eta), \quad \xi \in \Omega.$$

3 Green's Theorems on (Normal) Regions of the Sphere

A standard method for solving boundary value problems corresponding to the Beltrami operator Δ^* is the theory of Green's functions. We first introduce the definition and the properties of Green's function with respect to the Beltrami operator Δ^* (see [3]). Then we prove the fundamental theorem for L^* on normal regions on the sphere in order to solve differential problems involving the surface curl gradient (cf. [3, 6]).

Definition 3.1 The function $G(\Delta^*; \cdot, \cdot) : (\xi, \eta) \mapsto G(\Delta^*; \xi, \eta)$, $\xi, \eta \in \Omega$ with $-1 \leq \xi \cdot \eta < 1$, is called *Green's function* on Ω with respect to the operator Δ^* , if it satisfies the following properties:

1. (differential equation) For every point $\xi \in \Omega$, $\eta \mapsto G(\Delta^*; \xi, \eta)$ is twice continuously differentiable on $\{\eta \in \Omega : -1 \leq \xi \cdot \eta < 1\}$, and we have

$$\Delta_\eta^* G(\Delta^*; \xi, \eta) = -\frac{1}{4\pi}, \quad -1 \leq \xi \cdot \eta < 1.$$

2. (characteristic singularity) For every $\xi \in \Omega$, the function

$$\eta \mapsto G(\Delta^*; \xi, \eta) - \frac{1}{4\pi} \ln(1 - \xi \cdot \eta)$$

is continuously differentiable on Ω .

3. (rotational symmetry) For all orthogonal transformations \mathbf{t}

$$G(\Delta^*; \mathbf{t}\xi, \mathbf{t}\eta) = G(\Delta^*; \xi, \eta).$$

4. (normalization) For every $\xi \in \Omega$,

$$\int_{\Omega} G(\Delta^*; \xi, \eta) d\omega(\eta) = 0.$$

Following [3] the uniqueness of Green's function with respect to Δ^* is guaranteed. In fact, the function

$$G(\Delta^*; \xi, \eta) = \frac{1}{4\pi} \ln(1 - \xi \cdot \eta) + \frac{1}{4\pi} - \frac{1}{4\pi} \ln 2, \quad -1 \leq \xi \cdot \eta < 1,$$

is an explicit representation of Green's function with respect to the Beltrami operator Δ^* . In connection with (2.1) we obtain

$$L_\eta^* G(\Delta^*; \xi, \eta) = -\frac{1}{4\pi} \frac{\eta \wedge \xi}{1 - \xi \cdot \eta}, \quad -1 \leq \xi \cdot \eta < 1. \quad (3.1)$$

Next, we explain some geometrical assumptions imposed on subsets of the unit sphere $\Omega \subset \mathbb{R}^3$.

Definition 3.2 A region, i.e., an open and connected set $\Gamma \subset \Omega$, is called *normal* if the surface theorem of Stokes

$$\int_{\Gamma} L_\xi^* \cdot f(\xi) d\omega(\xi) = \int_{\partial\Gamma} \tau_\xi \cdot f(\xi) d\sigma(\xi)$$

is valid for all continuously differentiable vector fields $f \in c^{(1)}(\Omega)$, where $\partial\Gamma$ is the boundary curve of Γ , τ is the unit surface field pointing tangential to $\partial\Gamma$ and σ is the arc length along $\partial\Gamma$. A normal region $\Gamma \subset \Omega$ is called *regular*, if its boundary $\partial\Gamma$ has a continuously differentiable unit normal field $\nu : \partial\Gamma \rightarrow \mathbb{R}^3$ pointing outward of Γ , i.e., into $\Omega \setminus \bar{\Gamma}$.

By choosing $f = FL^*Q$, $F \in C^{(1)}(\bar{\Gamma})$, $Q \in C^{(2)}(\bar{\Gamma})$ in the surface theorem of Stokes we get Green's surface identity for the operator L^* , viz.

$$\int_{\Gamma} (L_{\eta}^*F(\eta) \cdot L_{\eta}^*Q(\eta) + F(\eta)\Delta_{\eta}^*Q(\eta)) \, d\omega(\eta) = \int_{\partial\Gamma} F(\eta)\tau_{\eta} \cdot L_{\eta}^*Q(\eta) \, d\sigma(\eta). \quad (3.2)$$

Let the function $F : \bar{\Gamma} \rightarrow \mathbb{R}$ be continuously differentiable and $\xi \in \bar{\Gamma}$ be fixed. Applying Green's surface identity to F and Green's function $G(\Delta^*; \xi, \cdot)$ on the region $\{\eta \in \partial\Gamma : |\xi - \eta| \geq \varepsilon\}$ we obtain for sufficiently small $\varepsilon > 0$

$$\begin{aligned} & \int_{|\xi-\eta| \geq \varepsilon, \eta \in \Gamma} (L_{\eta}^*F(\eta) \cdot L_{\eta}^*G(\Delta^*; \xi, \eta) + F(\eta)\Delta_{\eta}^*G(\Delta^*; \xi, \eta)) \, d\omega(\eta) \\ &= \int_{|\xi-\eta| = \varepsilon, \eta \in \Gamma} F(\eta)\tau_{\eta} \cdot L_{\eta}^*G(\Delta^*; \xi, \eta) \, d\sigma(\eta) + \int_{|\xi-\eta| \geq \varepsilon, \eta \in \partial\Gamma} F(\eta)\tau_{\eta} \cdot L_{\eta}^*G(\Delta^*; \xi, \eta) \, d\sigma(\eta), \end{aligned} \quad (3.3)$$

where σ denotes the arc length along $\partial\Gamma$ and $\{\eta \in \Gamma : |\xi - \eta| = \varepsilon\}$, while τ is the unit surface vector tangential to $\{\eta \in \bar{\Gamma} : |\xi - \eta| = \varepsilon\}$ or $\{\eta \in \partial\Gamma : |\xi - \eta| \geq \varepsilon\}$, respectively. Using Property 1 of Definition 3.1

$$\Delta_{\eta}^*G(\Delta^*; \xi, \eta) = -\frac{1}{4\pi}, \quad \eta \in \Gamma,$$

equation (3.3) can be rewritten as follows

$$\begin{aligned} & \int_{|\xi-\eta| \geq \varepsilon, \eta \in \Gamma} L_{\eta}^*F(\eta) \cdot L_{\eta}^*G(\Delta^*; \xi, \eta) \, d\omega(\eta) - \frac{1}{4\pi} \int_{|\xi-\eta| \geq \varepsilon, \eta \in \Gamma} F(\eta) \, d\omega(\eta) \\ &= \int_{|\xi-\eta| = \varepsilon, \eta \in \Gamma} F(\eta)\tau_{\eta} \cdot L_{\eta}^*G(\Delta^*; \xi, \eta) \, d\sigma(\eta) + \int_{|\xi-\eta| \geq \varepsilon, \eta \in \partial\Gamma} F(\eta)\tau_{\eta} \cdot L_{\eta}^*G(\Delta^*; \xi, \eta) \, d\sigma(\eta). \end{aligned} \quad (3.4)$$

Next, we concentrate on the integral

$$I_{\varepsilon}(\xi) = \int_{|\xi-\eta| = \varepsilon, \eta \in \Gamma} F(\eta)\tau_{\eta} \cdot L_{\eta}^*G(\Delta^*; \xi, \eta) \, d\sigma(\eta).$$

For each point $\eta \in \Gamma$ with $|\xi - \eta| = \varepsilon$, we have

$$\tau_{\eta} = \frac{\xi \wedge \eta}{\sqrt{1 - (\xi \cdot \eta)^2}}. \quad (3.5)$$

Hence, we find with (3.1)

$$I_{\varepsilon}(\xi) = -\frac{1}{4\pi} \int_{|\xi-\eta| = \varepsilon, \eta \in \Gamma} F(\eta) \frac{\sqrt{1 - (\xi \cdot \eta)^2}}{1 - \xi \cdot \eta} \, d\sigma(\eta).$$

Letting $\varepsilon \rightarrow 0$ we obtain, in analogy to well-known results of potential theory (see e.g. [3, 9]),

$$\lim_{\varepsilon \rightarrow 0} I_{\varepsilon}(\xi) = -\frac{\alpha(\xi)}{2\pi} F(\xi),$$

where $\alpha(\xi)$ is the solid angle subtended at $\xi \in \bar{\Gamma}$. This finally leads to the following integral formula.

Theorem 3.3 (Fundamental Theorem for L^* on Normal Regions) *Let Γ be a normal region with boundary $\partial\Gamma$. Suppose that F is a continuously differentiable function on $\bar{\Gamma}$, i.e., $F \in C^{(1)}(\bar{\Gamma})$. Then, for every point $\xi \in \Omega$, we have*

$$\frac{\alpha(\xi)}{2\pi} F(\xi) = \frac{1}{4\pi} \int_{\Gamma} F(\eta) \, d\omega(\eta) - \int_{\Gamma} L_{\eta}^*F(\eta) \cdot L_{\eta}^*G(\Delta^*; \xi, \eta) \, d\omega(\eta) + \int_{\partial\Gamma} F(\eta)\tau_{\eta} \cdot L_{\eta}^*G(\Delta^*; \xi, \eta) \, d\sigma(\eta).$$

Setting, particularly, $F = 1$ on $\bar{\Gamma}$ we immediately get from Theorem 3.3

$$\alpha(\xi) = \frac{\|\Gamma\|}{2} + 2\pi \int_{\partial\Gamma} \tau_{\eta} \cdot L_{\eta}^*G(\Delta^*; \xi, \eta) \, d\sigma(\eta), \quad \|\Gamma\| = \int_{\Gamma} d\omega. \quad (3.6)$$

Clearly, in case of a regular region, $\alpha(\xi) = 2\pi$ for all $\xi \in \Gamma$ and $\alpha(\xi) = \pi$ for all $\xi \in \partial\Gamma$. Furthermore, for the whole sphere Ω we have the following result (cf. [3, 4]).

Corollary 3.4 (Fundamental Theorem for L^* on Ω) *Suppose that F is of class $C^{(1)}(\Omega)$. Then, for every $\xi \in \Omega$,*

$$F(\xi) = \frac{1}{4\pi} \int_{\Omega} F(\eta) \, d\omega(\eta) - \int_{\Omega} L_{\eta}^*F(\eta) \cdot L_{\eta}^*G(\Delta^*; \xi, \eta) \, d\omega(\eta).$$

4 Potential and Stream Functions

Let us consider a continuous spherical vector field f of class $c^{(0)}(\Omega)$. For all $\xi \in \Omega$ we call $\xi \rightarrow f_{nor}(\xi) = (f(\xi) \cdot \xi)\xi$ the *normal field* of f , while $\xi \rightarrow f_{tan}(\xi) = f - f_{nor}(\xi)$, is called the *tangential field* of f . Obviously, $f(\xi) = f_{nor}(\xi) + f_{tan}(\xi)$ and $f_{nor}(\xi) \cdot f_{tan}(\xi) = 0$. Furthermore, for $f, g \in c^{(0)}(\Omega)$ and $\xi \in \Omega$, $f(\xi) \cdot g(\xi) = f_{nor}(\xi) \cdot g_{nor}(\xi) + f_{tan}(\xi) \cdot g_{tan}(\xi)$.

Lemma 4.1 *The tangential field of f vanishes, i.e., $f_{tan}(\xi) = 0$, $\xi \in \Omega$, if and only if $f(\xi) \cdot \hat{\tau}(\xi) = 0$ for every unit vector $\hat{\tau}(\xi)$ that is perpendicular to ξ , i.e., for which $\xi \cdot \hat{\tau}(\xi) = 0$, $\xi \in \Omega$.*

Proof. First, assume $f_{tan} = 0$. Then for all $\xi \in \Omega$, we have $f(\xi) \cdot \hat{\tau}(\xi) = 0$. Conversely, assume that the tangential field is non-vanishing, i.e., $f_{tan}(\xi) = f(\xi) - (f(\xi) \cdot \xi)\xi \neq 0$. Then it follows that $f_{tan}(\xi)|f_{tan}(\xi)|^{-1}$ is a unit vector field perpendicular to ξ . Hence, by our hypothesis, $f_{tan}(\xi) \cdot f_{tan}(\xi)|f_{tan}(\xi)|^{-1} = 0$. This implies $|f_{tan}(\xi)| = 0$ which is a contradiction. Thus it follows that $f_{tan}(\xi) = 0$, as required. \square

Lemma 4.2 *Suppose that f is continuous on a simply connected normal region $\Gamma \subset \Omega$. Moreover, let*

$$\int_{\mathcal{C}} \tau_{\xi} \cdot f(\xi) d\sigma(\xi) = 0$$

for every curve \mathcal{C} on Γ . Then $f_{tan}(\xi) = 0$ for all $\xi \in \Gamma$, i.e., the tangential field of f vanishes for all $\xi \in \Gamma$.

Proof. Choose any point $\xi_0 \in \Gamma$. Let τ_{ξ_0} be any unit vector satisfying $\tau_{\xi_0} \cdot \xi_0 = 0$. Then there is a curve \mathcal{C} on Γ passing through ξ_0 whose unit tangent vector at ξ_0 is just τ_{ξ_0} . Let $\mathcal{C}_{sub}^{\xi_0}$ be any subset of \mathcal{C} containing ξ_0 . Then, in accordance with our assumption,

$$\int_{\mathcal{C}_{sub}^{\xi_0}} \tau_{\xi} \cdot f(\xi) d\sigma(\xi) = 0.$$

Hence, letting the length of $\mathcal{C}_{sub}^{\xi_0}$ tend to zero we find $\tau_{\xi_0} \cdot f(\xi_0) = 0$. Lemma 4.1 then yields $f_{tan}(\xi_0) = f(\xi_0) - (f(\xi_0) \cdot \xi_0)\xi_0 = 0$. Since ξ_0 can be any point on Γ , we have $f_{tan}(\xi) = f(\xi) - (f(\xi) \cdot \xi)\xi = 0$ for all $\xi \in \Gamma$. This is the desired result. \square

The surface gradient acts like an ordinary gradient in \mathbb{R}^3 when we integrate it along lines on Γ . In more detail, suppose F is continuously differentiable in an open set in \mathbb{R}^3 containing Γ , and \mathcal{C} is any curve lying on Γ , starting at ξ_0 and ending at ξ_1 . Suppose that τ_{ξ} is the unit tangent vector at ξ on \mathcal{C} pointing from ξ_0 to ξ_1 . Then

$$F(\xi_1) - F(\xi_0) = \int_{\mathcal{C}} \tau_{\xi} \cdot \nabla_{\xi}^* F(\xi) d\sigma(\xi) \quad (4.1)$$

(observe that $\tau_{\xi} \cdot \nabla_{\xi} F(\xi) = \tau_{\xi} \cdot \nabla_{\xi}^* F(\xi)$, $\xi \in \Gamma$, see, e.g., [2, 4]). This result enables us to show the following lemma.

Lemma 4.3 *Suppose that Γ is a simply connected normal region. Let F be of class $C^{(1)}(\Gamma)$, then $\nabla_{\xi}^* F(\xi) = 0$ if and only if F is constant.*

Proof. If $\nabla_{\xi}^* F(\xi) = 0$, then we obtain, in connection with (4.1), $F(\xi_1) = F(\xi_0)$ for any ξ_0, ξ_1 on Γ . Conversely, if F is constant, the identity (4.1) shows that $f = \nabla^* F$ fulfills

$$\int_{\mathcal{C}} \tau_{\xi} \cdot f(\xi) d\sigma(\xi) = 0$$

for every curve \mathcal{C} lying on Γ . Consequently, following Lemma 4.2, $f_{tan}(\xi) = 0$ for all $\xi \in \Gamma$. This shows that $f_{tan}(\xi) = f(\xi) - (f(\xi) \cdot \xi)\xi = f(\xi) = \nabla_{\xi}^* F(\xi) = 0$ for all $\xi \in \Gamma$. \square

Next we prove the following result of spherical vector analysis (see, e.g., [2]).

Lemma 4.4 *Let $f \in c^{(0)}(\Gamma)$ be a tangent vector field on a simply connected region Γ , i.e., $f(\xi) = f_{tan}(\xi)$, $\xi \in \Gamma$. Furthermore, suppose that*

$$\int_{\mathcal{C}} \tau_{\xi} \cdot f(\xi) d\sigma(\xi) = 0$$

for every closed curve on Γ . Then there is a scalar field P on Γ such that

$$f(\xi) = \nabla_{\xi}^* P(\xi), \quad \xi \in \Gamma.$$

The field P is continuously differentiable and is unique up to a constant.

Proof. Take an arbitrary, but fixed $\xi_0 \in \Gamma$. We let

$$P(\xi) = \int_{\xi_0}^{\xi} \tau_{\zeta} \cdot f(\zeta) d\sigma(\zeta),$$

be the integral along any curve \mathcal{C} that starts at $\xi_0 \in \Gamma$ and ends at $\xi \in \Gamma$. Then, for any two points ξ_0, ξ on Γ and any curve \mathcal{C} lying on Γ and starting at ξ_0 and ending at ξ_1 ,

$$P(\xi_1) - P(\xi_0) = \int_{\xi_0}^{\xi_1} \tau_{\zeta} \cdot f(\zeta) d\sigma(\zeta). \quad (4.2)$$

Observing (4.1) we find

$$P(\xi_1) - P(\xi_0) = \int_{\xi_0}^{\xi_1} \tau_{\zeta} \cdot \nabla_{\zeta}^* P(\zeta) d\sigma(\zeta). \quad (4.3)$$

Combining (4.2) and (4.3) we obtain

$$\int_{\xi_0}^{\xi_1} \tau_{\zeta} \cdot (f(\zeta) - \nabla_{\zeta}^* P(\zeta)) d\sigma(\zeta) = 0$$

for any curve \mathcal{C} on Γ . Lemma 4.2, therefore, shows us that $f(\xi) - \nabla_{\xi}^* P(\xi) = 0$, $\xi \in \Gamma$. The proof that P is continuously differentiable on Γ is omitted. The easiest way to construct such a proof is to take P constant on each straight line passing through Γ in the normal direction (see, e.g., [2]). In order to verify that P is unique up to a constant, we observe that $\nabla_{\xi}^* P_1(\xi) = \nabla_{\xi}^* P_2(\xi)$, $\xi \in \Gamma$, implies $\nabla_{\xi}^*(P_1 - P_2)(\xi) = 0$, $\xi \in \Gamma$, i.e., by virtue of Lemma 4.3, $P_1 - P_2 = \text{const}$. \square

Now we are able to verify the following important theorem.

Theorem 4.5 *Let $f \in c^{(1)}(\Gamma)$ be a tangential field on a simply connected normal region Γ , i.e., $f(\xi) = f_{\tan}(\xi)$ for all $\xi \in \Gamma$. Then $L_{\xi}^* \cdot f(\xi) = 0$, $\xi \in \Gamma$, if and only if there is a scalar field P such that*

$$f(\xi) = \nabla_{\xi}^* P(\xi), \quad \xi \in \Gamma,$$

and P is unique up to an additive constant (P is called potential function for f).

Similarly, $\nabla_{\xi}^* \cdot f(\xi) = 0$, $\xi \in \Gamma$, if and only if there is a scalar field S such that

$$f(\xi) = L_{\xi}^* S(\xi), \quad \xi \in \Gamma,$$

and S is unique up to an additive constant (S is called stream function for f).

Proof. The condition $f = \nabla^* P$ implies $L^* \cdot f = 0$, and $f = L^* S$ implies $\nabla^* \cdot f = 0$.

Conversely, assume that $L_{\xi}^* \cdot f(\xi) = 0$, $\xi \in \Gamma$. Then the surface theorem of Stokes implies

$$\int_{\mathcal{C}} \tau_{\xi} \cdot f(\xi) d\sigma(\xi) = 0$$

for every closed curve \mathcal{C} on Γ . From Lemma 4.4 it follows that there exists a scalar field P such that $f = \nabla^* P$. Furthermore, P is unique up to an additive constant.

Finally, suppose $\nabla^* \cdot f = 0$. Then $L_{\xi}^* \cdot (\xi \wedge f(\xi)) = 0$, $\xi \in \Gamma$, hence, there is a scalar field S , unique up to a constant, such that $-\xi \wedge f(\xi) = \nabla_{\xi}^* S(\xi)$, $\xi \in \Gamma$. This is equivalent to $-\xi \wedge (\xi \wedge f(\xi)) = (\xi \wedge \nabla_{\xi}^*) S(\xi)$, $\xi \in \Gamma$, or $f = L^* S$ on Γ . This proves Theorem 4.5. \square

From Lemma 4.3 we are immediately able to deduce the following statement.

Lemma 4.6 *Let F be of class $C^{(1)}(\Gamma)$, then $L_{\xi}^* F(\xi) = 0$ if and only if F is constant.*

Proof. If $L_{\xi}^* F(\xi) = 0$, i.e., $\xi \wedge \nabla_{\xi}^* F(\xi) = 0$ for all $\xi \in \Gamma$. Then $\xi \wedge \xi \wedge \nabla_{\xi}^* F(\xi) = \xi \cdot \nabla_{\xi}^* F(\xi) - \nabla_{\xi}^* F(\xi)(\xi \cdot \xi) = -\nabla_{\xi}^* F(\xi) = 0$ for all $\xi \in \Gamma$. Thus by virtue of Lemma 4.3 we find $F = \text{const}$.

Conversely, if F is constant, then $L_{\xi}^* F(\xi) = \xi \wedge \nabla_{\xi}^* F(\xi) = 0$ for all $\xi \in \Gamma$. This proves our assertion. \square

5 The Differential Equations of the Surface Curl Gradient

In what follows we give two different solvability conditions for the differential equation (1.1) of the surface curl gradient L^* . First, based on the results of Chapter 3 we formulate a certain integrability condition to assure uniqueness.

Theorem 5.1 *Given $f = L^*F \in c^{(0)}(\Omega)$. Then the scalar function F is uniquely determined by the condition:*

$$\frac{1}{4\pi} \int_{\Omega} F(\eta) d\omega(\eta) = C_0, \quad C_0 \in \mathbb{R}.$$

Proof. We suppose that $F_1, F_2 \in C^{(1)}(\Omega)$ are functions satisfying the condition above. Then the difference $D = F_1 - F_2$ satisfies $L^*D = 0$ on Ω and $\frac{1}{4\pi} \int_{\Omega} D(\eta) d\omega(\eta) = 0$. Consequently, from Corollary 3.4, we obtain $D(\xi) = 0$ for all $\xi \in \Omega$. Therefore, $F_1 = F_2$, as required. \square

Theorem 5.2 *Let $\Gamma \subset \Omega$ be a simply connected normal region. Given $f = L^*F \in c^{(0)}(\Omega)$. Then the scalar function F is uniquely determined by the condition taken at one point $\xi_0 \in \bar{\Gamma}$:*

$$\frac{1}{4\pi} \int_{\Gamma} F(\eta) d\omega(\eta) + \int_{\partial\Gamma} F(\eta) \tau_{\eta} \cdot G(\Delta^*; \xi_0, \eta) d\sigma(\eta) = C_0, \quad C_0 \in \mathbb{R}.$$

Proof. We look at the difference D of two solutions which satisfies $L^*D = 0$ in Γ . By Lemma 4.6 we find $D(\xi) = \text{const} = C$ for all $\xi \in \bar{\Gamma}$. In connection with Theorem 3.3 we have

$$\frac{C}{2\pi} \left(\frac{\|\Gamma\|}{2} + 2\pi \int_{\partial\Gamma} \tau_{\eta} \cdot L_{\eta}^* G(\Delta^*; \xi_0, \eta) d\sigma(\eta) \right) = 0.$$

Using with (3.6) we, therefore, find $C = 0$, i.e., $D = 0$ on $\bar{\Gamma}$, as required. \square

Second, based on the results of Chapter 4, we are able to formulate a uniqueness condition by fixing a certain functional value.

Theorem 5.3 *Given $f = L^*F \in c^{(0)}(\Omega)$. Then the scalar function F is uniquely determined by the condition taken at one point $\xi_0 \in \Omega$:*

$$F(\xi_0) = C_0, \quad C_0 \in \mathbb{R}.$$

Proof. The constant difference D of two functions satisfying the conditions is equal to $D(\xi_0) = C_0 - C_0 = 0 = D(\xi)$ for all $\xi \in \Omega$. \square

Theorem 5.4 *Suppose that Γ is a simply connected normal region. Given $f = L^*F \in c^{(0)}(\Omega)$. Then the scalar function F is uniquely determined by the condition taken at one point $\xi_0 \in \bar{\Gamma}$:*

$$F(\xi_0) = C_0, \quad C_0 \in \mathbb{R}.$$

Proof. D is constant on $\bar{\Gamma}$ with $D(\xi_0) = 0$. Hence $D(\xi) = 0$ in $\bar{\Gamma}$. \square

6 Regularized Green's Theorems on (Normal) Regions on the Sphere

In the following we first introduce the regularized Green function with respect to Δ^* . We state its definition together with some properties to be needed for the discussion of spherical wavelets on regular regions.

Definition 6.1 Given $\rho \in (0, 2)$, the regularized Green function with respect to Δ^* is defined for all $\xi, \eta \in \Omega$ by

$$G_{\rho}(\Delta^*; \xi, \eta) = \begin{cases} \frac{1}{4\pi} \ln(1 - \xi \cdot \eta) + \frac{1}{4\pi} - \frac{1}{4\pi} \ln 2, & 1 - \xi \cdot \eta > \rho, \\ \frac{1}{4\pi\rho} (1 - \xi \cdot \eta) + \frac{1}{4\pi} \ln(\rho) - \frac{1}{4\pi} \ln 2, & 1 - \xi \cdot \eta \leq \rho. \end{cases}$$

The regularized Green function with respect to the Beltrami operator $(\xi, \eta) \rightarrow G_{\rho}(\Delta^*; \xi, \eta)$ only depends on the inner product of ξ and η , hence, it is a radial basis function, i.e., $G_{\rho}(\Delta^*; \mathbf{t}\xi, \mathbf{t}\eta) = G_{\rho}(\Delta^*; \xi, \eta)$ holds true for all orthogonal transformations \mathbf{t} . Figure 1 gives an illustration of the regularized Green function with respect to Δ^* . Note that, by construction, this kernel function represents an approximation of the original Green's function, i.e., it converges pointwise to Green's function as ρ tends to 0.

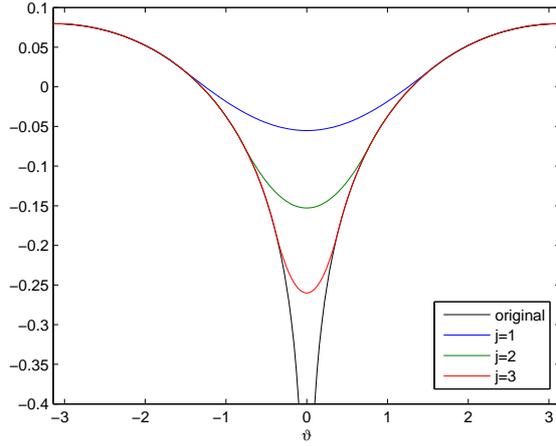


Fig. 1 The regularized Green function $\vartheta \mapsto G_\rho(\Delta^*; \cos \vartheta)$ for $\rho = 1 - \cos(\pi 2^{-j})$ with $j = 1, 2, 3$ and the original Green function $G_\rho(\Delta^*; \xi \cdot \eta)$ with respect to the Beltrami operator Δ^* . Note that $\xi \cdot \eta = \cos \vartheta$, $\vartheta = \angle(\xi, \eta)$.

We immediately realize that the regularized Green function with respect to Δ^* is continuously differentiable. Applying the surface curl gradient L^* to the second variable yields to the so-called *regularized Green function with respect to L^** . Obviously, for $\rho \in (0, 2)$, we obtain for all $\xi, \eta \in \Omega$

$$g_\rho^{L^*}(\xi, \eta) = L_\eta^* G_\rho(\Delta^*; \xi, \eta) = \begin{cases} \frac{1}{4\pi} \frac{1}{1-\xi \cdot \eta} (\xi \wedge \eta), & 1 - \xi \cdot \eta > \rho, \\ \frac{1}{4\pi} \frac{1}{\rho} (\xi \wedge \eta), & 1 - \xi \cdot \eta \leq \rho. \end{cases} \quad (6.1)$$

Observing the equation $|\xi \wedge \eta| = \sqrt{1 - (\xi \cdot \eta)^2}$ we derive for all $\xi, \eta \in \Omega$ and $\rho \in (0, 2)$

$$|g_\rho^{L^*}(\xi, \eta)| = \begin{cases} \frac{1}{4\pi} \sqrt{\frac{1+\xi \cdot \eta}{1-\xi \cdot \eta}}, & 1 - \xi \cdot \eta > \rho, \\ \frac{1}{4\pi} \sqrt{1 - (\xi \cdot \eta)^2}, & 1 - \xi \cdot \eta \leq \rho. \end{cases}$$

A graphical impression of the norm of the regularized Green function with respect to L^* and the norm of the surface curl gradient of Green's function with respect to Δ^* is illustrated in Figure 2. By similar arguments as known from potential theory (see e.g. [6, 9]) we obtain the following counterpart of the integral formula developed in Chapter 3.

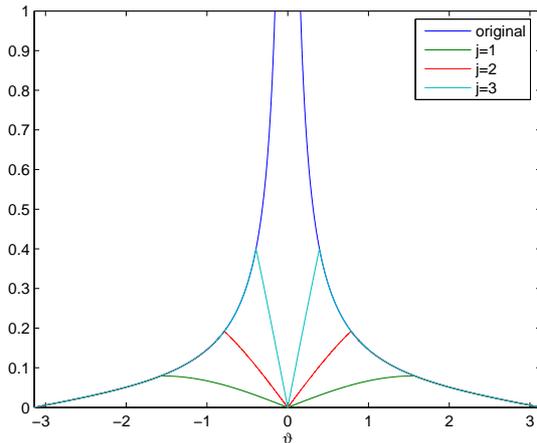


Fig. 2 The norm of the regularized Green function $\vartheta \mapsto |g_\rho^{L^*}(\cos(\vartheta))|$ for $\rho = 1 - \cos(\pi 2^{-j})$ with $j = 1, 2, 3$ and the norm of the surface curl gradient of Green's function with respect to Δ^* .

Theorem 6.2 For $F \in C^{(1)}(\Omega)$ we have

$$\limsup_{\rho \rightarrow 0} \limsup_{\xi \in \Omega} \left| F(\xi) - \frac{1}{4\pi} \int_{\Omega} F(\eta) d\omega(\eta) - \int_{\Omega} g_\rho^{L^*}(\xi, \eta) \cdot L^* F(\eta) d\omega(\eta) \right| = 0.$$

After deriving the regularized version of the integral theorem for L^* on Ω we now turn to the regularized integral theorem for L^* on normal regions Γ . For that purpose we introduce the following settings.

Definition 6.3 Let $\Gamma \subset \Omega$ be a normal region of the unit sphere Ω . For $F \in C^{(1)}(\bar{\Gamma})$ we let

$$S_\rho(F)(\xi) = \int_\Gamma g_\rho^{L^*}(\xi, \eta) \cdot L_\eta^* F(\eta) d\omega(\eta) - \int_{\partial\Gamma} g_\rho^{L^*}(\xi, \eta) \cdot F(\eta) \tau_\eta d\sigma(\eta), \quad \rho \in (0, 2),$$

as a counterpart of

$$S(F)(\xi) = \int_\Gamma L_\eta^* G(\Delta^*; \xi, \eta) \cdot L_\eta^* F(\eta) d\omega(\eta) - \int_{\partial\Gamma} L_\eta^* G(\Delta^*; \xi, \eta) \cdot F(\eta) \tau_\eta d\sigma(\eta).$$

Clearly, it is not hard to show that

$$\limsup_{\rho \rightarrow 0} \sup_{\xi \in \Omega} |S(F)(\xi) - S_\rho(F)(\xi)| = 0. \quad (6.2)$$

Theorem 6.4 (Regularized Integral Theorem for L^* on Γ) Let $\Gamma \subset \Omega$ be a normal region with boundary $\partial\Gamma$. Suppose that F is a continuously differentiable function on $\bar{\Gamma}$, i.e., $F \in C^{(1)}(\bar{\Gamma})$. Then

$$\limsup_{\rho \rightarrow 0} \sup_{\xi \in \bar{\Gamma}} \left| \frac{\alpha(\xi)}{2\pi} F(\xi) - \frac{1}{4\pi} \int_\Gamma F(\eta) d\omega(\eta) + S_\rho(F)(\xi) \right| = 0,$$

where $\alpha(\xi)$ denotes, as usually, the solid angle subtended at $\xi \in \bar{\Gamma}$.

7 Vector Spherical Wavelets on Normal Regions

We turn our attention to the introduction of vector spherical wavelets, where $\Gamma \subset \Omega$ is supposed to be a normal region. We choose a sequence which divides the continuous scale interval $(0, 2)$ into discrete pieces. More explicitly, $(\rho_j)_{j \in \mathbb{N}_0}$ denotes a sequence of real numbers satisfying $\lim_{j \rightarrow \infty} \rho_j = 0$ and $\lim_{j \rightarrow 0} \rho_j = 1$. For example, we can choose $\rho_j = 2^{-j}$ or $\rho_j = 1 - \cos(\pi 2^{-j})$, $j \in \mathbb{N}_0$. The point of departure for our considerations on normal regions Γ is Theorem 6.4 in the form

$$\frac{\alpha(\xi)}{2\pi} F(\xi) - F_{mean}^\Gamma = - \lim_{j \rightarrow \infty} S_{\rho_j}(F)(\xi), \quad \text{where } F_{mean}^\Gamma = \frac{1}{4\pi} \int_\Gamma F(\eta) d\omega(\eta), \quad \xi \in \bar{\Gamma}. \quad (7.1)$$

Note that the discrete steps in this approximation process are called *scales*, i.e., the value j takes the role of the scale parameter, i.e., the "zooming-in" parameter. By using discrete regularization parameters we are naturally led to the following type of scale discretized Green wavelets.

Definition 7.1 Let $\{g_{\rho_j}^{L^*}\}_{j \in \mathbb{N}_0}$ be the regularized Green function with respect to L^* (see (6.1)). Then the scale discretized regularized Green wavelet function with respect to L^* is defined by

$$\psi_{\rho_j} = g_{\rho_{j+1}}^{L^*} - g_{\rho_j}^{L^*} \quad (7.2)$$

$$= L_\eta^* G_{\rho_{j+1}}(\Delta^*; \xi, \eta) - L_\eta^* G_{\rho_j}(\Delta^*; \xi, \eta), \quad j \in \mathbb{N}_0. \quad (7.3)$$

In fact, the difference of two consecutive scales of regularized Green functions with respect to Δ^* reads

$$\begin{aligned} & G_{\rho_{j+1}}(\Delta^*; \xi, \eta) - G_{\rho_j}(\Delta^*; \xi, \eta) \\ &= \begin{cases} 0, & 1 - \xi \cdot \eta > \rho_j, \\ \frac{1}{4\pi} \ln(1 - \xi \cdot \eta) - \frac{1}{4\pi\rho_j} (1 - 1 - \xi \cdot \eta) + \frac{1}{4\pi} (1 - \ln(\rho_j)), & \rho_j > 1 - \xi \cdot \eta > \rho_{j+1}, \\ \left(\frac{1}{4\pi\rho_{j+1}} - \frac{1}{4\pi\rho_j} \right) (1 - \xi \cdot \eta) + \frac{1}{4\pi} (\ln(\rho_{j+1}) - \ln(\rho_j)), & 1 - \xi \cdot \eta \leq \rho_{j+1}, \end{cases} \end{aligned}$$

such that

$$\psi_{\rho_j}(\xi, \eta) = g_{\rho_{j+1}}^{L^*}(\xi, \eta) - g_{\rho_j}^{L^*}(\xi, \eta) = \begin{cases} 0, & 1 - \xi \cdot \eta > \rho_j, \\ \frac{1}{4\pi} \left(\frac{1}{\rho_j} - \frac{1}{1 - \xi \cdot \eta} \right) (\xi \wedge \eta), & \rho_j > 1 - \xi \cdot \eta > \rho_{j+1}, \\ \frac{1}{4\pi} \left(\frac{1}{\rho_j} - \frac{1}{\rho_{j+1}} \right) (\xi \wedge \eta), & 1 - \xi \cdot \eta \leq \rho_{j+1}. \end{cases}$$

A graph of the norm of the scale discretized regularized Green wavelet function with respect to L^* for the discretization parameters $\rho = 1 - \cos(\pi 2^{-j})$ with $j = 0, 1, 2, 3$ is shown in Figure 3. Note, that the functions ψ_{ρ_j} have a local support.

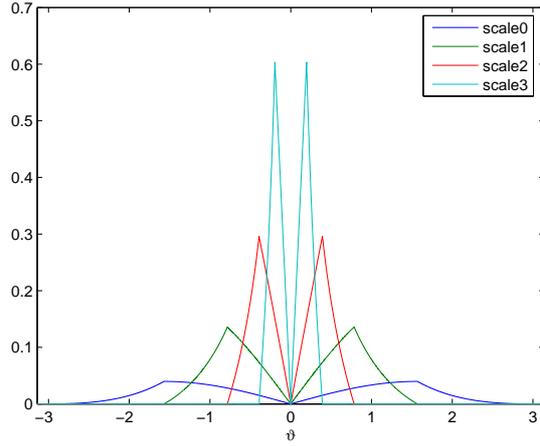


Fig. 3 The norm of the regularized Green wavelet function $\vartheta \rightarrow |\Psi_{\rho_j}(\cos(\vartheta))|$ with respect to L^* for $\rho = 1 - \cos(\pi 2^{-j})$ with scale $j = 0, 1, 2, 3$.

$S_{\rho_j}(F)(\xi)$, as given by Definition 6.3, is called the *scale discrete regularized Green scaling function transform* with respect to L^* . Let $\{\psi_{\rho_j}\}_{j \in \mathbb{N}_0}$ be the scale discretized regularized Green function with respect to L^* . The scale discretized regularized Green wavelet transform with respect to L^* is defined by

$$W_{\rho_j}(F)(\xi) = \int_{\Gamma} L_{\eta}^* F(\eta) \cdot \psi_{\rho_j}(\xi, \eta) d\omega(\eta) - \int_{\partial\Gamma} F(\eta) \tau_{\eta} \cdot \psi_{\rho_j}(\xi, \eta) d\sigma(\eta).$$

We arrive at the following theorem.

Theorem 7.2 Let $\{g_{\rho_j}^{L^*}\}_{j \in \mathbb{N}_0}$ be the regularized Green function with respect to L^* . Then the multiscale reconstruction of a function $F \in C^{(1)}(\Gamma)$ is given by

$$\frac{\alpha(\xi)}{2\pi} F(\xi) - F_{mean}^{\Gamma} = \sum_{j=-\infty}^{\infty} W_{\rho_j}(F)(\xi), \quad \xi \in \Gamma,$$

where the equality holds in the $\|\cdot\|_{C(\bar{\Gamma})}$ -sense.

By observing the definition of the scaling transform $S_{\rho_j}(F)(\xi)$, Theorem 7.2 admits the following reformulation.

Corollary 7.3 Under the assumptions of Theorem 7.2

$$\frac{\alpha(\xi)}{2\pi} F(\xi) - F_{mean}^{\Gamma} - S_{\rho_j}(F)(\xi) = \sum_{j=J}^{\infty} W_{\rho_j}(F)(\xi), \quad \xi \in \bar{\Gamma},$$

for every $J \in \mathbb{N}_0$ in the $\|\cdot\|_{C(\bar{\Gamma})}$ -sense.

These reconstruction formula will now be applied to the modelling of oceanic circulation.

8 Modelling of Sea Surface Topography From Ocean Circulation

As is well known, the (spherical) Earth Ω can be decomposed uniquely into two normal regions, namely the oceanic part Ω^{oc} and the continental part Ω^{co} . Due to the definition of the sea surface topography we assume the differential equation $L^*H = h$ to be valid on Ω^{oc} thereby relating the sea surface topography $H \in C^{(1)}(\bar{\Omega}^{oc})$ to the geostrophic ocean flow $h \in C^{(0)}(\bar{\Gamma})$. Moreover, we assume that $H = 0$ on $\partial\Omega^{oc}$ and $\int_{\Omega^{oc}} H(\xi) d\omega(\xi) = 0$. Summarizing our results for the normal region $\Gamma = \Omega^{oc}$, we therefore obtain

$$H(\xi) = - \int_{\Omega^{oc}} L_{\eta}^* G(\Delta^*; \xi, \eta) \cdot h(\eta) d\omega(\eta), \quad \xi \in \Omega^{oc},$$

which can be approximated as follows

$$H(\xi) = - \lim_{j \rightarrow \infty} \int_{\Omega^{oc}} L_{\eta}^* G_{\rho_j}(\Delta^*; \xi, \eta) \cdot h(\eta) d\omega(\eta), \quad \xi \in \Omega^{oc}$$

It suffices to have an (in the sense of Weyl) equidistributed nodal set $(\eta_i, h(\eta_i))$, $\eta_i \in \overline{\Omega^{oc}}$, $i = 1, \dots, N$, to discretize the integral on the left hand side and to establish a multiscale approximation of the sea surface topography.

In practice, however, we are often confronted with the problem of determining the sea surface topography on a certain subdomain Γ of Ω^{oc} (e.g., caps, squares or rectangles), where suitable discrete data information about the velocity vectors of ocean currents is available. When we are interested in solving that problem numerically from discrete data material our approach shows that we have to know, in addition, the sea surface topography on the boundary $\partial\Gamma$. But this quantity does not vanish generally. Even more, our numerical calculation based on discrete data is only unique up to a constant H_{mean}^Γ :

$$\frac{\alpha(\xi)}{2\pi} H(\xi) - H_{mean}^\Gamma = - \lim_{j \rightarrow \infty} \left(\int_{\Gamma} h(\eta) \cdot g_{\rho_j}^{L^*}(\xi, \eta) d\omega(\eta) + \int_{\partial\Gamma} H(\eta) \tau_\eta \cdot g_{\rho_j}^{L^*}(\xi, \eta) d\sigma(\eta) \right).$$

In what follows, particular attention is paid to the numerical stability caused by the specific observation of the boundary terms in our calculation. Two examples will be presented for the discrete "Topography Problem" involving discrete data on the boundary $\partial\Gamma$. First, we consider a spherical cap Γ as a regular region. In a second step we will have a look at a normal region, i.e., a rectangle.

8.1 Multiscale Approximation on a Spherical Cap

The spherical cap under consideration is defined by its center $\zeta \in \Omega$ and its radius $r > 0$, more precisely we let

$$\Gamma_r(\zeta) = \{\eta \in \Omega : |\zeta - \eta| < r\}.$$

In this case the tangential unit vector τ_η is explicitly given for all $\eta \in \partial\Gamma_r^{sea}(\zeta)$ by (3.5). Substituting the tangential unit vector in the equations above we obtain

$$\frac{\alpha(\xi)}{2\pi} H(\xi) - H_{mean}^{\Gamma_r} = - \lim_{j \rightarrow \infty} \left(\int_{\Gamma_r(\zeta)} h(\eta) \cdot g_{\rho_j}^{L^*}(\xi, \eta) d\omega(\eta) + \int_{\partial\Gamma_r(\zeta)} H(\eta) \frac{g_{\rho_j}^{L^*}(\xi, \eta) \cdot (\zeta \wedge \eta)}{\sqrt{1 - (\zeta \cdot \eta)^2}} d\sigma(\eta) \right).$$

To be more precise, the region of interest in our first example is a spherical cap Γ_{21} where 21° denotes the apex angle of the cap. Furthermore, we assume that the velocity field measurements are not continuously given, but on an equiangular longitude-latitude grid with a step size of 0.1° . The sea surface topography is prescribed at a finite set of boundary points that are sampled with an angular distance of 0.0072° . Both data sets have been generated from an artificial topography model, the CLS01 model (which the Geomatics Group, TU Kaiserslautern, received from the French enterprise CLS (Collecte Localisation Satellites)). It has been computed using a 7-year TOPEX/ POSEIDON mean profile, a 5-year ERS-1/2 mean profile, a 2-year GEOSAT mean profile and the two 168-day non repeat cycle data of the ERS-1 geodetic phase. The data were processed and homogenized using the most recent corrections to compute the CLS01 mean sea surface (CLS01 MSS). The surface is derived from altimetric data on oceans and from geoidal undulations elsewhere (continuously connected in between, starting from an ocean depth of 10 meters). Since the velocity as well as the sea surface topography have been generated from the CLS01 model we are able to compare the calculated results with the original data later on. Figure 4 illustrates the mean dynamic topography which is used to calculate the input dataset, i.e., the geostrophic flow which is also shown in Figure 4. Since we are especially interested in boundary effects, we always plot the spherical cap together with its surrounding environment.

We deal with the discretization of the scale interval $(0, 2)$. Very often the regularization parameter is set to $\rho_j = 2^{-j}$, $j \in \mathbb{N}_0$. In order to increase the speed of convergence we use the definition $\rho_j = 1 - \cos(2^{-2j})$, $j \in \mathbb{N}_0$, instead. The effect is that we obtain a good approximation of the sea surface topography already at low scales, which can be seen in Figure 5 presenting the reconstruction of the topography at lower scales. The left column contains the reconstruction based on the regularized Green function with respect to L^* and the right column shows the approximation of the topography with the corresponding scale discretized regularized Green wavelet function.

At first sight the approximated sea surface topography is close to the original topography inside the spherical cap, while the topography at the boundary of the spherical cap seems to be disturbed. Taking the absolute value of the difference between the original and the approximated sea surface topography, we can better specify the approximation errors (see Figure 6).

8.2 Multiscale Approximation on a Spherical Rectangle

Next, we will have a look at the aforementioned rectangular region. From the dataset of the spherical cap we cut off a rectangular region and generate a corresponding boundary dataset. Figure 7 illustrates the dynamic topography and the geostrophic flow which is used as input dataset.

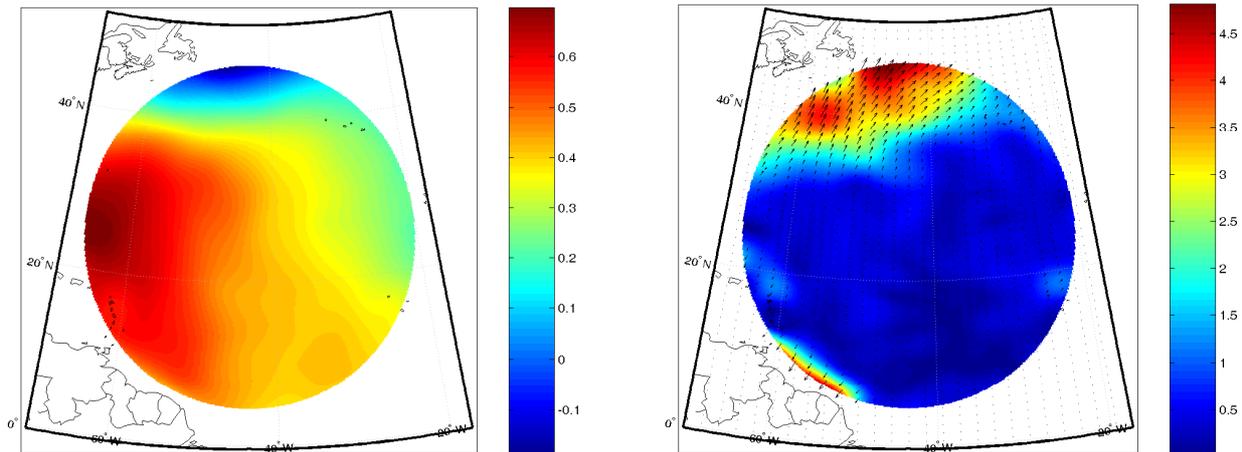


Fig. 4 Plot of the mean dynamic topography in [m] (left), which is used to calculate the input dataset, i.e., the geostrophic flow in [cm/s] (right).

Omitting the vertices in the process of integration it is easy to calculate τ_η for each $\eta \in \partial\Gamma$. As in the case of the spherical cap we achieve a good approximation of the sea surface topography, which can be seen in Figure 8. It shows the reconstruction of the topography at lower scales. The left column contains the reconstruction based on the regularized Green function with respect to L^* and the right column shows the approximation of the topography with the corresponding regularized Green wavelet function. Analogously to the case of a spherical cap we achieve small errors.

A question we are interested in is in how far the part of the boundary integral influences the reconstruction. To give an impression of this influence Figure 10 compares the total reconstruction to the part of the reconstruction coming from the boundary integral in the cap and in the rectangular region at scale 5.

It is an interesting fact that the boundary integral takes the major part of the whole reconstruction. The numerical results demonstrate that the new multiscale approximation method based on the regularized Green function with respect to L^* yields an efficient procedure to model sea surface topography from ocean velocity field data. Already at low scales, i.e., with larger regularization caps, the approximation error inside of the regular region tends to zero. Compared to standard methods using a truncated velocity field model that neglects the boundary effects our approximation method shows an appealing convergence of the approximated version to the exact solution at the boundary of the region.

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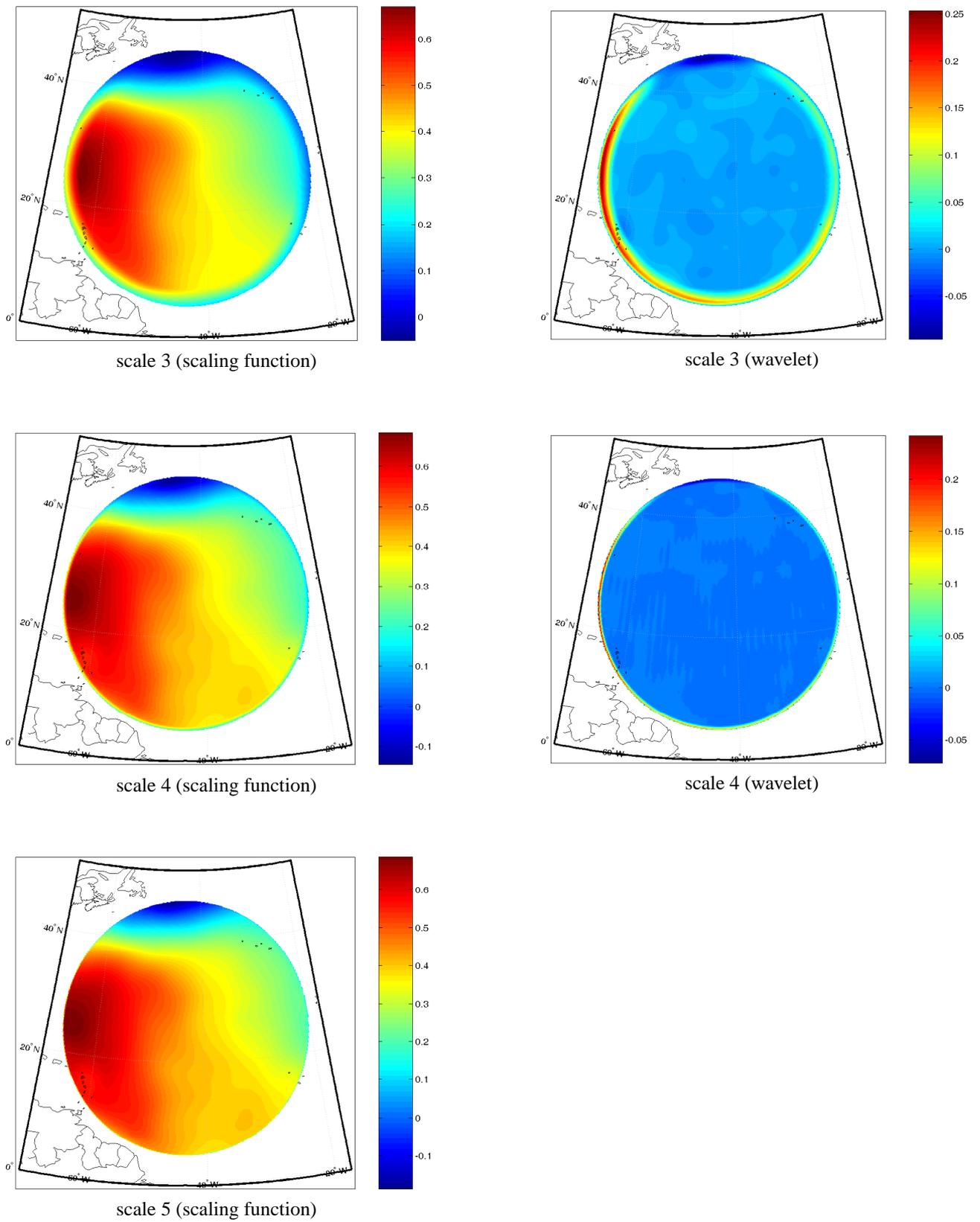


Fig. 5 Reconstruction of the topography in [m] on the spherical cap Γ_{21} at low scales using the scale discretized regularized Green function.

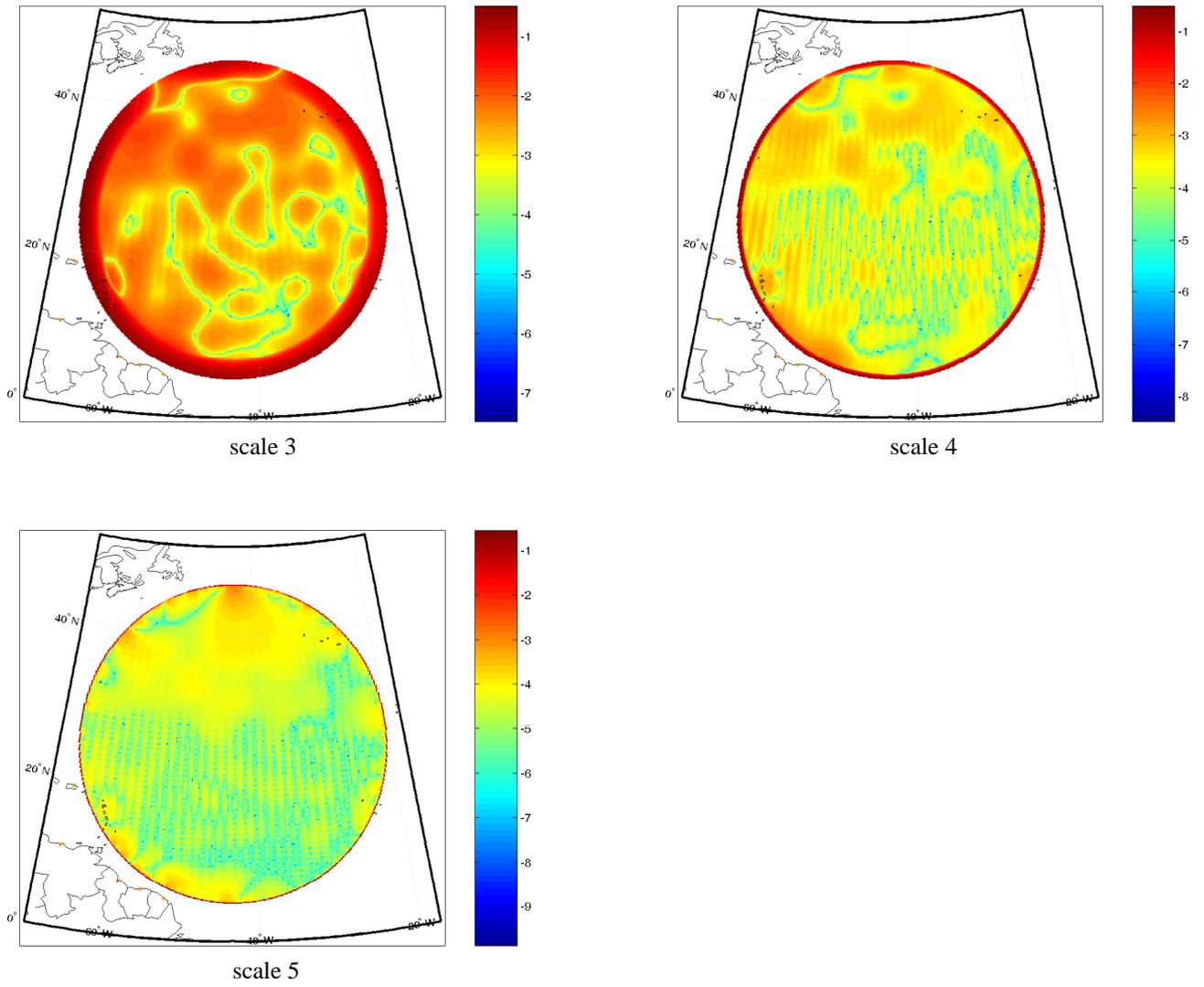


Fig. 6 Approximation error of the dynamic topography in the spherical cap Γ_{21} , where the constant value H_{mean}^Γ was subtracted. Note that the colorbar is logarithmic.

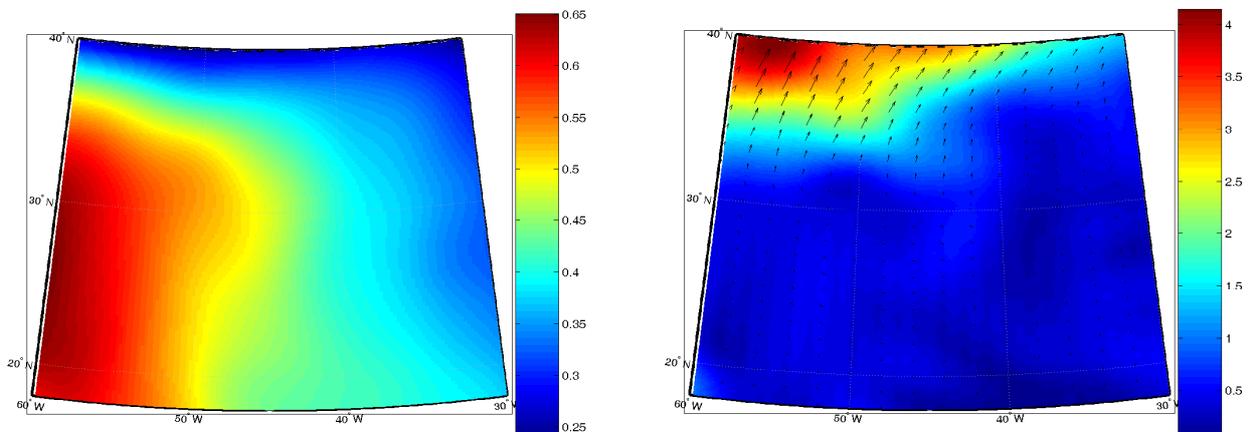


Fig. 7 Plot of the mean dynamic topography in [m] (left) and the geostrophic flow in [cm/s] (right) in a rectangular region .

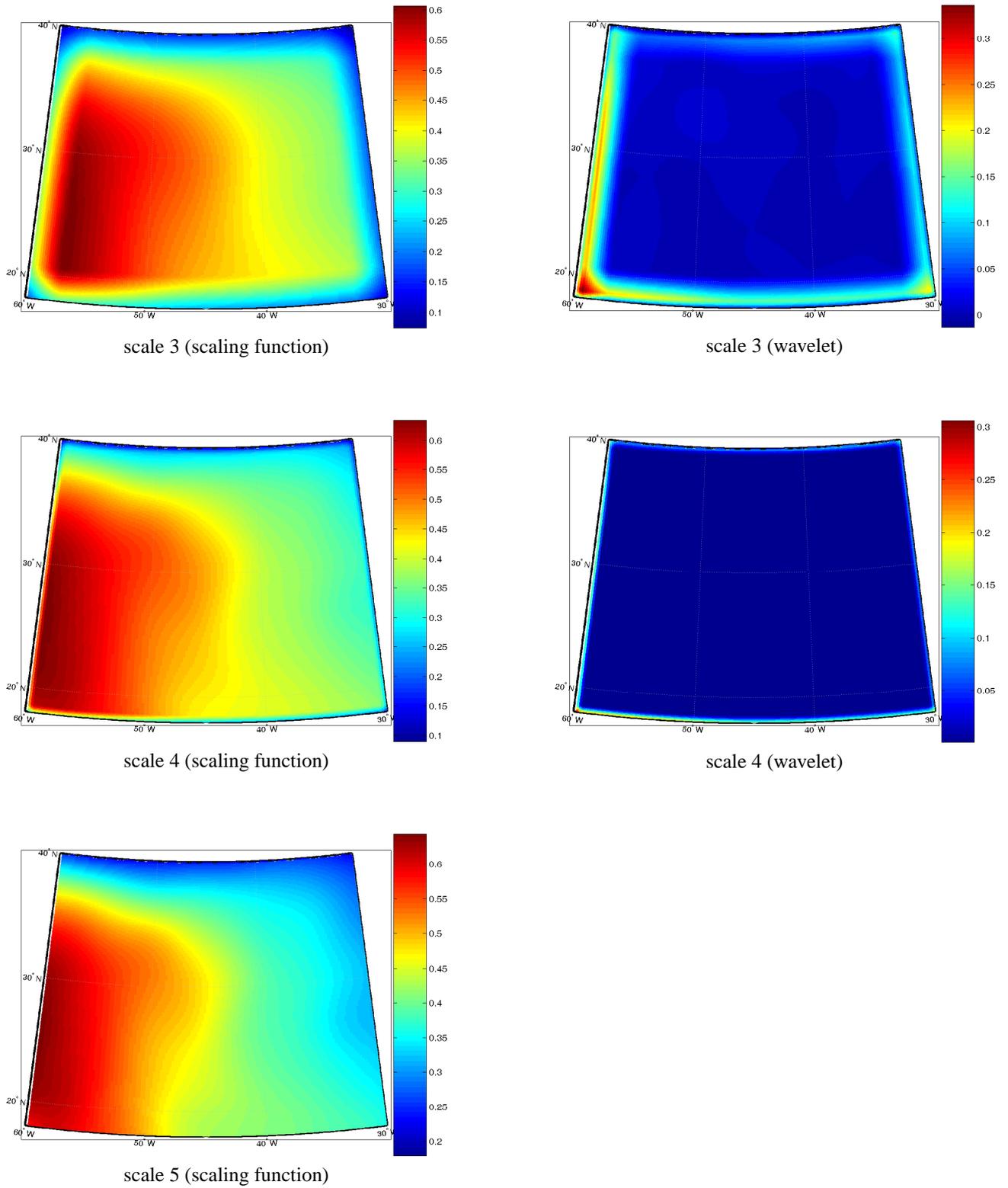


Fig. 8 Reconstruction of the topography in [m] on a rectangular region Γ at low scales using the scale discretized regularized Green function.

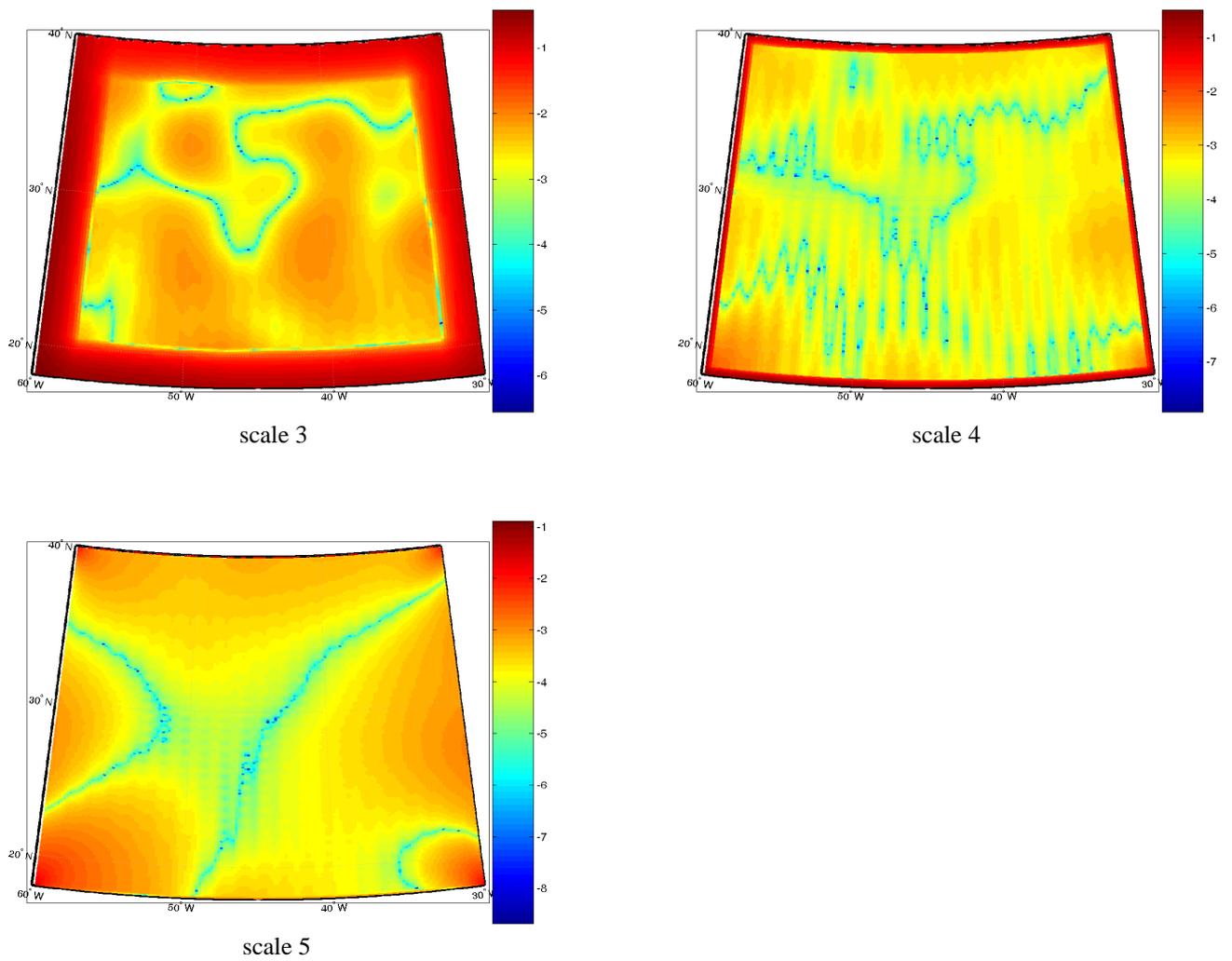


Fig. 9 Approximation error of topography in a rectangular region Γ , where the constant value H_{mean}^Γ was subtracted. Note that the colorbar is logarithmic.

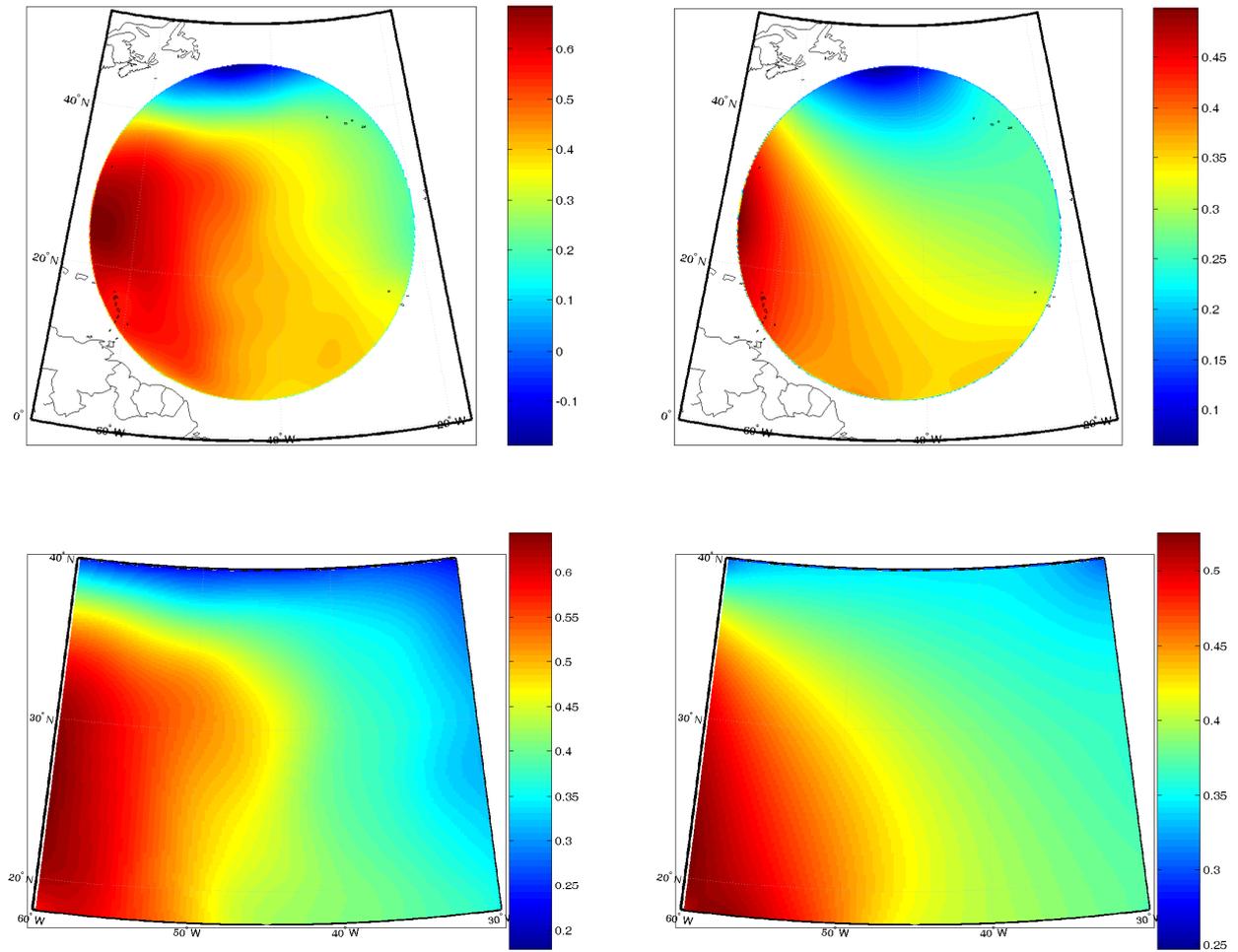


Fig. 10 Comparison of the total reconstruction of the dynamic topography in [m] (left) to the part of the reconstruction coming from the boundary integral (right) in the cap and in the rectangular region at scale 5.

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KAISERSLAUTERN

Informationen:

Prof. Dr. W. Freeden

Prof. Dr. E. Schock

Fachbereich Mathematik

Technische Universität Kaiserslautern

Postfach 3049

D-67653 Kaiserslautern

E-Mail: freeden@mathematik.uni-kl.de

schock@mathematik.uni-kl.de